

RESONANCE OF A DRIVEN OSCILLATOR

PHYS 2302, Saint Mary's University

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In class, we found the amplitude and phase of a damped harmonic oscillator driven by a force $F(t) = F_0 \cos \omega t$ are given by:

$$A(\omega) = \frac{F_0/m}{\sqrt{(\omega_0^2 - \omega^2)^2 + 4\gamma^2\omega^2}}; \quad (1)$$

$$\phi(\omega) = \tan^{-1} \frac{2\gamma\omega}{\omega_0^2 - \omega^2}. \quad (2)$$

where: m = mass of oscillator;

ω = driving frequency;

$\omega_0 = \sqrt{\frac{k}{m}}$ = oscillation frequency of undamped system;

$\gamma = \frac{b}{2m}$ = damping coefficient.

Also shown in class, $A(\omega)$ is a maximum at the *resonant frequency*,

$$\omega_r = \sqrt{\omega_0^2 - 2\gamma^2},$$

which is real so long as $\gamma < \frac{\omega_0}{\sqrt{2}}$. For $\omega = \omega_r$, Eq. (1) becomes:

$$A_{\max} = \frac{F_0}{m\sqrt{\omega_0^4 - \omega_r^4}}. \quad (3)$$

Before plotting equations (1) – (3), we first scale them. Let $\alpha = \gamma/\omega_0$ and $\beta = \omega/\omega_0$, both unitless. Then,

$$A = \frac{F_0}{\underbrace{m\omega_0^2}_k \sqrt{(1 - \beta^2)^2 + 4\alpha^2\beta^2}}$$

$$\Rightarrow \tilde{A}(\beta) \equiv \frac{Ak}{F_0} = \frac{1}{\sqrt{(1 - \beta^2)^2 + 4\alpha^2\beta^2}}, \quad (4)$$

where \tilde{A} is also a unitless quantity.

Next,

$$\tan \phi(\beta) = \frac{2\gamma\omega}{\omega_0^2 - \omega^2} \frac{\omega_0^2}{\omega_0^2} = \frac{2\alpha\beta}{1 - \beta^2}. \quad (5)$$

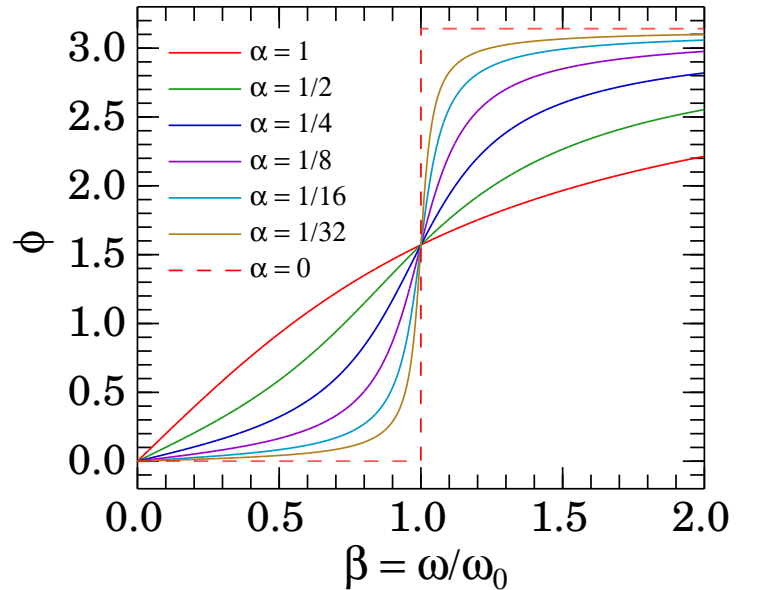
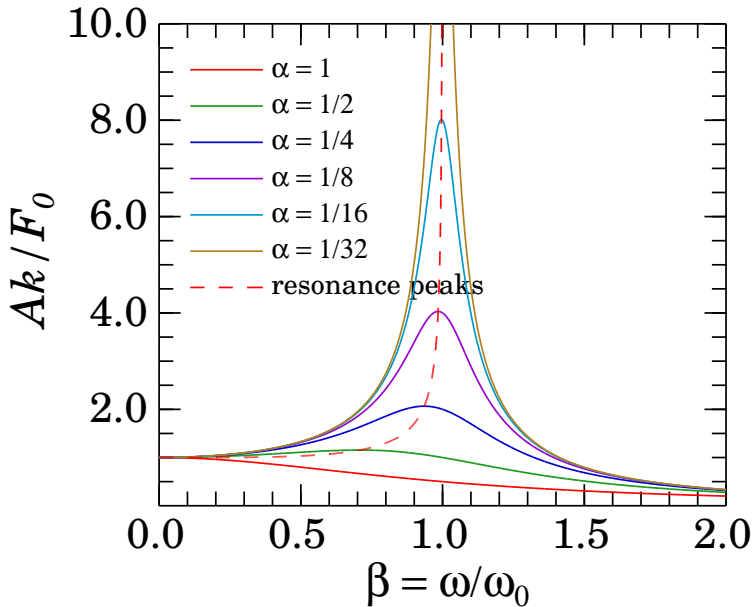
Finally, for $\beta_r \equiv \omega_r/\omega_0$,

$$A_{\max} = \frac{F_0}{m\omega_0^2\sqrt{1 - \beta_r^4}} \Rightarrow \tilde{A}_{\max} = \frac{A_{\max}k}{F_0} = \frac{1}{\sqrt{1 - \beta_r^4}}, \quad (6)$$

for $0 < \beta_r < 1$.

Equations (4) – (6) can now be plotted by specifying:

- a *domain* for β [e.g., $\in (0 : 2)$, where $\beta = 1 \Rightarrow$ system is driven at the undamped oscillation frequency];
- α , unitless damping coefficient, γ/ω_0 ;
- β , unitless driving frequency, ω/ω_0 .



Left: Unitless amplitude, \tilde{A} , plotted against unitless driving frequency, β , for various values of unitless damping coefficient, α [Eq. (3)]. The peak amplitude for each α occurs at the unitless resonance frequency,

$$\beta_r = \frac{\omega_r}{\omega_0} = \frac{1}{\omega_0} \sqrt{\omega_0^2 - 2\gamma^2} = \sqrt{1 - 2\alpha^2}.$$

The dashed red line is \tilde{A}_{\max} plotted against $\beta = \beta_r \in (0, 1)$ [Eq. (6)]. Since $\beta_r \in \mathbb{R}$, there can be no resonant frequency for $\alpha > \frac{1}{\sqrt{2}} \sim 0.707$. Thus, there is an A_{\max} (barely) for $\alpha = 0.5$ (green), but not for $\alpha = 1$ (solid red) where the plot decreases monotonically from $\beta = 0$.

Right: Phase lag, ϕ , plotted against β for various α [Eq. (5)]. As $\alpha \rightarrow 0$ (no damping), $\phi(\beta)$ becomes a step function (dashed red line), with $x(t)$ and $F(t)$ in phase ($\phi = 0$) for $\omega < \omega_0$, and of opposite phase ($\phi = \pi$) for $\omega > \omega_0$.

A lesson on the importance of understanding resonance: *Gallopig Gertie*, the Tacoma Narrows Bridge built in 1940.

