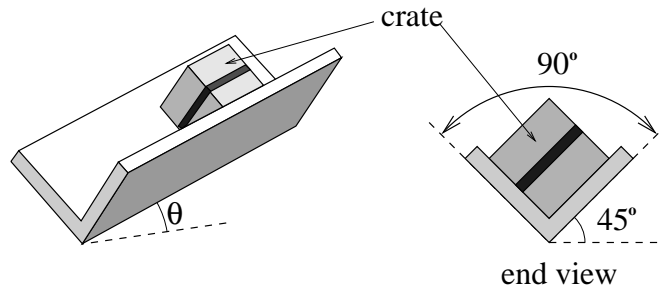


First Practise Midterm Solutions

PHYS 2302 (Mechanics I); D. A. Clarke

Problem 1. Crate sliding down a V-ramp: A crate slides down a right-angle trough inclined to the horizontal at an angle θ as shown in the figure. If the coefficient of kinetic friction between the sides of the crate and the surface of the trough is μ_k , find the acceleration of the crate.

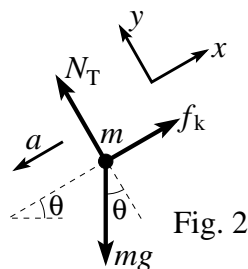
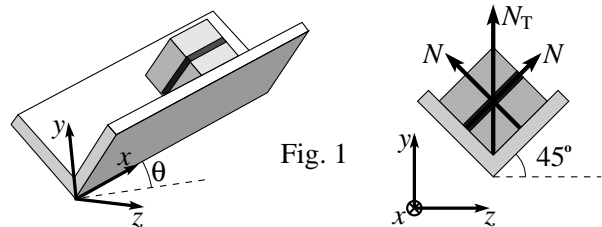


Hint: Normal forces are exerted on the crate by both sides of the trough.

Solution: From Fig. 1 to the right, we see that the z -components of the normal forces cancel, leaving a total normal force $\vec{N}_T = \sqrt{2}N\hat{j}$. This simplification allows us to construct our free body diagram (Fig. 2 below) entirely in the x - y plane.

On the other hand, the kinetic frictional force responds to the normal forces from each surface directly, and $\vec{f}_k = \mu_k(2N)\hat{i}$.

Thus, from the FBD, Newton's 2nd Law for this problem is written:



$$x/ \quad f_k - mg \sin \theta = 2\mu_k N - mg \sin \theta = -ma; \quad (1)$$

$$y/ \quad N_T - mg \cos \theta = \sqrt{2}N - mg \cos \theta = 0 \quad (2)$$

From equation (2), $N = \frac{mg}{\sqrt{2}} \cos \theta$. Substituting this into equation (1), we get:

$$2\mu_k \frac{mg}{\sqrt{2}} \cos \theta - mg \sin \theta = -ma \quad \Rightarrow \quad \boxed{a = g(\sin \theta - \sqrt{2}\mu_k \cos \theta)}.$$

Short and sweet.

Problem 2. Sliding block: A metal block slides along a horizontal surface lubricated with a heavy oil that imposes on the block a resistive force given by:

$$F(v) = -cv^{3/2}, \quad (3)$$

where v is the speed of the block, c is a constant, and the negative sign indicates this force causes the block to slow down.

- a) (12 points) If at $x = 0$ the speed of the block is v_0 , show that the speed of the block after it slides a distance x is given by,

$$v(x) = \left(\sqrt{v_0} - \frac{cx}{2m} \right)^2.$$

- b) (2 points) Find X , the maximum distance travelled by the block.
c) (6 points) By solving the differential equation,

$$F(v) = m \frac{dv}{dt},$$

determine whether the block reaches X asymptotically (*i.e.*, as $t \rightarrow \infty$) or after a finite time, T .

Solution: a) To find velocity as a function of distance, $v(x)$, we need Newton's second law in terms of v and x . Thus,

$$F = ma = m \frac{dv}{dt} = m \frac{dx}{dt} \frac{dv}{dx} = mv \frac{dv}{dx} = -cv^{3/2},$$

which is a separable first order ODE in v and x . Thus,

$$v^{-1/2} dv = -\frac{c}{m} dx \quad \Rightarrow \quad \int_{v_0}^v (v')^{-1/2} dv' = -\frac{c}{m} \int_0^x dx',$$

where integration is taken from $x = 0$ where the speed of the block is v_0 to a distance x where the speed is v . Doing the integrals, we get:

$$\begin{aligned} 2\sqrt{v'} \Big|_{v_0}^v &= -\frac{c}{m} x' \Big|_0^x \quad \Rightarrow \quad 2(\sqrt{v} - \sqrt{v_0}) = -\frac{cx}{m} \\ &\Rightarrow \quad \boxed{v(x) = \left(\sqrt{v_0} - \frac{cx}{2m} \right)^2}, \end{aligned} \quad (4)$$

as desired.

- b) To find the maximum distance, X , the block slides, set $v = 0$ in equation (4) to get:

$$0 = \left(\sqrt{v_0} - \frac{cX}{2m} \right)^2 \quad \Rightarrow \quad \boxed{X = \frac{2m\sqrt{v_0}}{c}}.$$

c) To find the time it takes for the block to stop, we start with Newton's second law in terms of v and t :

$$F = m \frac{dv}{dt} = -cv^{3/2} \Rightarrow v^{-3/2} dv = -\frac{c}{m} dt \Rightarrow \int_{v_0}^0 v^{-3/2} dv = -\frac{c}{m} \int_0^T dt$$

$$\Rightarrow -\frac{2}{\sqrt{v}} \Big|_{v_0}^0 = -\frac{c}{m} t \Big|_0^T \Rightarrow T = \frac{2m}{c} \frac{1}{\sqrt{v}} \Big|_{v_0}^0 \rightarrow \infty.$$

Thus, the block reaches its maximum distance asymptotically. This points to a weakness in the model for the viscosity of the oil, equation (3), since we know perfectly well that the block comes to a dead stop well within our lifetimes!

Problem 3. SHO plug and chug: Consider the position of a simple harmonic oscillator as a function of time given by,

$$x(t) = C \cos(\omega_0 t - \phi_0), \tag{5}$$

where C is the amplitude of oscillation, ω_0 is the angular frequency of oscillation, and ϕ_0 is the phase of oscillation determining where the oscillator is at $t = 0$.

- a) (5 points) Show by direct substitution that equation (5) solves the differential equation of motion for a simple harmonic oscillator, as given on the formula sheets.
- b) (15 points) If, at $t = 0$, $x(0) = x_0$ and $v(0) = \dot{x}(0) = v_0$, find C and ϕ_0 in terms of quantities given.

Solution: a) As given on the formula sheets, the differential equation of motion for an SHO is:

$$\frac{d^2 x}{dt^2} = -\omega_0^2 x. \tag{6}$$

Substituting equation (5) into the LHS of equation (6), we get:

$$\frac{d^2 x}{dt^2} = -\omega_0^2 C \cos(\omega_0 t - \phi_0) = -\omega_0^2 x,$$

as desired.

b) For this part, we need the velocity of the oscillator given by the time-derivative of equation (5):

$$v(t) = \frac{dx}{dt} = -C\omega_0 \sin(\omega_0 t - \phi_0). \tag{7}$$

Then, applying initial conditions, we get:

$$x(0) = C \cos(-\phi_0) = C \cos \phi_0 = x_0; \tag{8}$$

$$v(0) = -C\omega_0 \sin(-\phi_0) = v_0 \Rightarrow C \sin \phi_0 = \frac{v_0}{\omega_0}. \quad (9)$$

Dividing equation (9) by equation (8), we get:

$$\frac{C \sin \phi_0}{C \cos \phi_0} = \frac{v_0/\omega_0}{x_0} \Rightarrow \boxed{\tan \phi_0 = \frac{v_0}{\omega_0 x_0}}.$$

Next, adding the squares of equations (8) and (9), we get:

$$C^2 \cos^2 \phi_0 + C^2 \sin^2 \phi_0 = C^2 = x_0^2 + \frac{v_0^2}{\omega_0^2} \Rightarrow \boxed{C = \sqrt{x_0^2 + \frac{v_0^2}{\omega_0^2}}}.$$
